

磁场作用下Oldroyd-B流体的非稳态驻点流动

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摘要

本文研究了拉伸板上Oldroyd-B流体的非稳态驻点流动, 此外, 还考虑了磁场对流动的影响。利用合适的相似变换将控制方程转化为常微分方程。结合双参数变形展开方法与基函数方法(DPTEM-BF)求得微分方程的近似解析解。通过图形展示了各物性参数对速度场的影响, 值得注意的是: 随着松弛时间参数的增加, 流体的粘性力增大使得流体流动阻力增强, 这导致了流体流速减小; 延迟时间参数的变化使得速度场增大, 并且动量边界层厚度变厚; 较大的磁场参数对应于较大的洛伦兹力, 这个力阻碍了流体流动, 因而流体流动速度降低。

关键词

Oldroyd-B流体, 非稳态驻点流动, 磁场, DPTEM-BF方法

Unsteady Stagnation-Point Flow of Oldroyd-B Fluid along a Stretching Sheet with Magnetic Field

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Abstract

This paper examines the influence of magnetic field on unsteady stagnation-point flow of Oldroyd-B

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fluid towards a stretching sheet. Suitable similarity transformations are put into use to yield the ordinary differential equation, which are dealt with double-parameter transformation expansion method with base function method (DPTEM-BF). Impacts of various physical parameters on the velocity field are explored via graphs. It is noteworthy that as relaxation time parameter enlarges, the viscous force of fluid increases, which causes larger resistance to fluid flow. The increase of retardation time parameter increases both the velocity and momentum boundary layer thickness. Larger magnetic parameter corresponds to the larger Lorentz force, which impedes fluid flow and thus slows it down.

Keywords

Oldroyd-B Fluid, Unsteady Stagnation-Point Flow, Magnetic Field, DPTEM-BF Method

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1. 引言

近年来,由于非牛顿流体在工业、食品以及日常生活中的广泛应用而备受研究人员的关注,如石油、油漆、淀粉液、蛋清、果浆等都是非牛顿流体。其中, Oldroyd-B 流体就是常见的一种非牛顿流体,它不仅能够反映流体的松弛特性,还能表征延迟特性对流体的影响。Khan 等[1]研究了 Oldroyd-B 流体流动问题。Hayat 等[2]对多孔拉伸板上 Oldroyd-B 流体的混合自然流动进行了讨论。Ullah 等[3]得到了两同轴圆柱之间 Oldroyd-B 流体旋转流动的精确解。此外, Alshomrani 等[4]还探索了在分层机制作用下,热辐射和化学反应对 Oldroyd-B 纳米流体流动和传热的影响。Hayat 等[5]在 Alshomrani 等人的基础之上还考虑了热产生/吸收的影响并且推导了其关系式。Hafeez 等[6]采用 Cattaneo-Christov 双扩散模型分析了多孔旋转圆盘上 Oldroyd-B 流体的流动、传热和传质,并且得到了数值解。

驻点流动在熔融纺丝工艺、拉丝、玻璃吹制、金属和聚合物挤出等应用中起着至关重要的作用,因此不少学者都对驻点流动问题进行了研究。Mustafa 等[7]研究了幂律流体的驻点流动。Sharma 等[8]对具有二阶速度滑移的拉伸板上微极流体的驻点流动进行了分析。在布朗运动和热泳扩散存在下, Bai 等[9]对拉伸板上 Maxwell 纳米磁流体的驻点流动也进行了讨论。在此基础之上, Bai 等[10]还采用 Cattaneo-Christov 双扩散模型探索了具有化学反应和热辐射效应的上随体 Oldroyd-B 纳米磁流体的驻点流动和传热。上述学者都只研究了稳态驻点流动,但实际上流体的流动与时间也有关,以下研究人员对不同流体的非稳态驻点流动也进行了分析。Hayat 等[11]讨论了拉伸板上二级流体的非稳态驻点流动问题。Zaib 等[12]研究了速度滑移和热泳效应对微极流体非稳态驻点流动的影响。Awais 等[13]还分析了在焦耳加热下,径向拉伸圆盘上 Maxwell 磁流体的非稳态驻点流动。Mahdy [14]对具有 Biot 数效应的加热拉伸板上 Oldroyd-B 纳米流体的非稳态驻点流动进行了建模,并且还考虑了纳米流体的动态粘度和导热率的影响。此外, Hayat 等[15]也对 Oldroyd-B 纳米流体的非稳态驻点流动问题进行了探索,并且将热产生/吸收、非线性热辐射、布朗运动以及热泳扩散的影响考虑在内。

磁流体由于同时具有独特的磁性以及流动性在扬声器、阻尼器和选矿分离等方面有着独特的价值,这一特性使得不少学者对存在磁场情况下的流体的流动问题进行了深入思考。Asghar 等[16]得到了在磁场作用下粘性流体的库埃特流动的解析解。Hayat 等[17]研究了倾斜磁场对指数拉伸板上三级流体流动的影响。Mabood 等[18]对在磁场和加热作用下,拉伸板上 Jeffery 流体的非稳态流动问题进行了讨论,并且还考虑了变导热系数的影响。Jakati 等[19]分析了在斜向外磁场作用下,非均匀热源/散热器、布朗运动和

热泳对线性拉伸板板上 Maxwell 纳米流体流动的影响。此外, Hashmi 等[20]还研究了磁场对两个无限等温拉伸盘间 Oldroyd-B 纳米流体的混合对流以及轴对称流动的影响。然而, 很少有学者研究了磁场对 Oldroyd-B 流体非稳态驻点流动的影响。

本文主要研究了在磁场作用下, 拉伸板上 Oldroyd-B 流体的非稳态驻点流动。采用合适的相似变换将偏微分方程转化为常微分方程, 并且使用双参数变形展开方法与基函数方法结合的方法(DPTEM-BF)进行求解。最后通过图形分析了松弛时间参数、延迟时间参数、磁场参数等对速度场的影响。

2. 数学模型

考虑拉伸板上 Oldroyd-B 流体的非稳态驻点流动问题。建立如图 1 所示的流动示意图, 其中, 拉伸板的方向沿着 x 轴的方向, y 轴垂直于拉伸板的方向。假设拉伸板的速度为 $U_w = bx/(1-at)$, 驻点流动速度为 $U_e = cx/(1-at)$, 其中 $a, b, c > 0$ 是常数。此外, 沿着 y 轴方向施加了一个强度为 B_0 的均匀磁场。

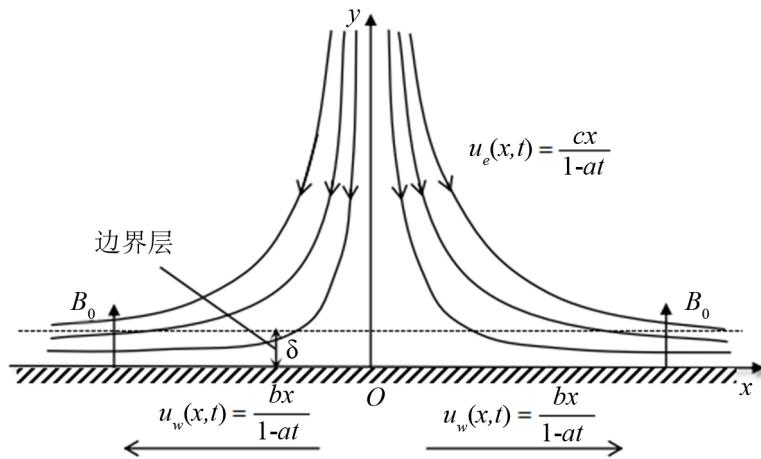


Figure 1. The physical flow diagram

图 1. 流动示意图

Oldroyd-B 流体的非稳态驻点流动控制方程为:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \quad (1)$$

$$\begin{aligned} \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = -\lambda_1 & \left(\frac{\partial^2 u}{\partial t^2} + 2u \frac{\partial^2 u}{\partial x \partial t} + 2v \frac{\partial^2 u}{\partial y \partial t} + u^2 \frac{\partial^2 u}{\partial x^2} + 2uv \frac{\partial^2 u}{\partial x \partial y} + v^2 \frac{\partial^2 u}{\partial y^2} \right) \\ & + v \frac{\partial^2 u}{\partial y^2} + \lambda_2 v & \left(\frac{\partial^3 u}{\partial y^2 \partial t} + v \frac{\partial^3 u}{\partial y^3} + u \frac{\partial^3 u}{\partial x \partial y^2} - \frac{\partial u}{\partial x} \frac{\partial^2 u}{\partial y^2} - \frac{\partial u}{\partial y} \frac{\partial^2 v}{\partial y^2} \right) \\ & + \frac{\partial u_e}{\partial t} + u_e \frac{\partial u_e}{\partial x} + \lambda_1 u_e^2 \frac{\partial^2 u_e}{\partial x^2} - \frac{\sigma B_0^2}{\rho} & \left[u - u_e + \lambda_1 \left(\frac{\partial u}{\partial t} + v \frac{\partial u}{\partial y} - \frac{\partial u_e}{\partial t} \right) \right] \end{aligned} \quad (2)$$

边界条件为:

$$\text{当 } y = 0 \text{ 时, } u = u_w(x, t) = \frac{bx}{1-at}, \quad v = 0 \quad (3)$$

$$\text{当 } y \rightarrow \infty \text{ 时, } u = u_e(x, t) = \frac{cx}{1-at} \quad (4)$$

其中, u 和 v 分别为沿 x 和 y 轴方向上的速度分量, $\nu = \mu/\rho$ 指的是流体的运动粘度系数, λ_1 和 λ_2 分别是

松弛时间参数和延迟时间参数。

采用如下的无量纲函数以及相似变量:

$$u = \frac{bx}{(1-at)} f'(\eta), v = -\sqrt{\frac{vb}{(1-at)}} f(\eta), \eta = \sqrt{\frac{b}{v(1-at)}} y, \psi = \sqrt{\frac{vb}{(1-at)}} xf(\eta) \quad (5)$$

将(5)带入(2)~(4), 得到对应的常微分方程及边界条件如下:

$$\begin{aligned} & S\left(f' + \frac{1}{2}f''\eta\right) + (f')^2 - ff'' + \beta_1 \left\{ S^2 \left(2f' + \frac{7}{4}f''\eta + \frac{1}{4}f'''^2\right) + 2S \left[(f')^2 + \frac{1}{2}ff''\eta\right] \right. \\ & \left. - S(3ff'' + f'''^2) - 2ff'' + f^2f'''\right\} - f''' - \beta_2 \left[S \left(2f'' + \frac{1}{2}f^{iv}\eta\right) - ff^{iv} + (f'')^2 \right] \\ & - AS - A^2 + M \left\{ (f' - A) + \beta_1 \left[S \left(f' + \frac{1}{2}f''\eta\right) - ff'' - SA \right] \right\} \end{aligned} \quad (6)$$

对应的边界条件转化为:

$$\text{当 } \eta = 0 \text{ 时, } f(0) = 0, \quad f'(0) = 1 \quad (7)$$

$$\text{当 } \eta \rightarrow \infty \text{ 时, } f'(\infty) = A \quad (8)$$

其中, $\beta_1 = \lambda_1 b / (1-at)$ 和 $\beta_2 = \lambda_2 b / (1-at)$ 分别为松弛时间参数和延迟时间参数, $S = a/b$ 为不稳定参数, $A = c/b$ 是驻点参数, $M = \sigma B_0^2 (1-at) / \rho b$ 指的是磁场参数。

3. DPTEM-BF 方法求解

采用双参数变形展开方法与基函数结合的方法求解微分方程, 根据双参数变形展开方法的原理, 对自变量和因变量进行变形, 引入小参数 ε 和新变量 ξ , 将函数 $f(\eta)$ 转化为:

$$f(\eta) = \varepsilon^4 F(\xi) + f(0) + f'(0)\eta + \frac{f''(0)}{2!}\eta^2 + \frac{f'''(0)}{3!}\eta^3, \quad \xi = \varepsilon^{-1}\eta \quad (9)$$

令 $f''(0) = a_1$, $f'''(0) = a_2$, 其中 a_1 , a_2 均为未知常数, $F(\xi)$ 为变量 ξ 的函数。

对应地, 函数 $F(\xi)$ 的初值条件为:

$$F(0) = 0, \quad F'(0) = 0, \quad F''(0) = 0, \quad F'''(0) = 0 \quad (10)$$

将 $F(\xi)$ 在 $\varepsilon = 0$ 附近展开成级数形式:

$$F(\xi) = F_0(\xi) + \sum_{i=1}^{\infty} F_i(\xi) \varepsilon^i \quad (11)$$

将(10)带入(11)中, 得到:

$$F_i(0) = 0, \quad F'_i(0) = 0, \quad F''_i(0) = 0, \quad F'''_i(0) = 0, \quad i = 0, 1, 2, 3, \dots \quad (12)$$

将(9)和(12)带入式(6)中, 把微分方程转化为关于 ε 的多项式, 根据多项式恒为零的性质, 使多项式中 ε 同次幂的系数为零, 再利用初值条件, 便可以计算出 $F_i(\xi)$ 的表达式, 紧接着便可以得到 $F(\xi)$ 的表达式, 消去小参数 ε , 得到含有 a_1 , a_2 的关于 $f(\eta)$ 的表达式。采用结合基函数来确定参数 a_1 , a_2 的值, 将函数 $f(\eta)$ 表示成如下基函数的组合:

$$f(\eta) \approx f_{N_1, N_2}(\eta) = f_{0,0}(\eta) + \sum_{j=3}^{N_1} \sum_{i=1}^{N_2} a_{ij} f_{i,j}(\eta) = f_{0,0}(\eta) + \sum_{j=3}^{N_1} \sum_{i=1}^{N_2} a_{ij} \eta^j e^{ia_0 \eta} \quad (13)$$

其中, $f_{0,0}(\eta)$ 为:

$$f_{0,0}(\eta) = m_1 + s\eta - m_1 e^{a_0\eta} + (1-s+a_0m_1)\eta e^{a_0\eta} + m_2 \eta^2 e^{a_0\eta} \quad (14)$$

令 $N_1 = N_2 = 3$ ，其中 $m_1, m_2, a_0, a_{13}, a_{23}, a_{33}$ 为假设的未知数，最后使用牛顿迭代法便可以计算出每个未知参数的值，将其带入函数 $f(\eta)$ 的表达式，就可以得到原微分方程的近似解析解。

4. 结果与讨论

采用双参数变形展开方法与基函数结合的方法求得了微分方程的近似解析解，由表 1 可知，将所求结果与已有的研究结果进行比较，具有很好的一致性。此外，通过图形详细分析讨论了各物性参数对速度场的影响。

Table 1. Comparison of $-f''(0)$ with different values of S while $\beta_1 = \beta_2 = A = M = 0$

表 1. 当 $\beta_1 = \beta_2 = A = M = 0$ 时， S 取不同值， $-f''(0)$ 值的比较结果

S	Sharidan et al. [21]	Chamkha et al. [22]	Mukhopadhyay et al. [23]	Present results
0.8	1.261042	1.261512	1.261479	1.261826
1.2	1.377722	1.378052	1.377850	1.377496

图 2 展示了松弛时间参数 β_1 对速度场的影响，由图 2 可知，当 β_1 增加时，流体的粘性力增大，这使得流体的流动阻力增大，因此流体的流动速度减小。

图 3 绘制了不同延迟参数值 β_2 下的速度分布情况，从图中可以发现，延迟时间参数对速度场有着相反的影响，随着 β_2 的增大，流体动量穿透深度加深，拉伸板上流体的速度分量增加，这意味着流体流速增加。此外， β_2 的增大也会使得流体动量边界层厚度变厚。

图 4 反映了不稳定参数 S 变化时速度场的变化趋势，由于 S 的增加，流体的流动速度以及边界层厚度都均匀地减小。

图 5 给出了驻点参数 A 对速度的影响，当 A 增加时，流体的动量损失减小，边界层内流体速度越快趋于周围流体速度，因此速度边界层厚度会变薄。

图 6 显示了磁场参数 M 变化时对速度的影响，随着 M 值的增大，流体的流动速度减慢，并且边界层厚度变薄。从物理上来讲， M 的增加使得磁场作用下的洛伦兹力增大，这个力阻碍了流体流动，因此流体流动速度减小。

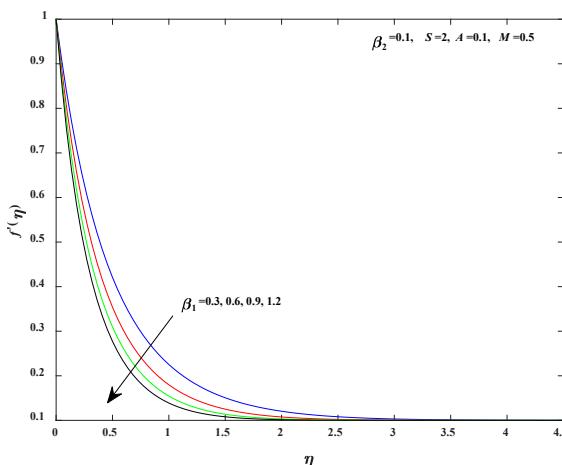
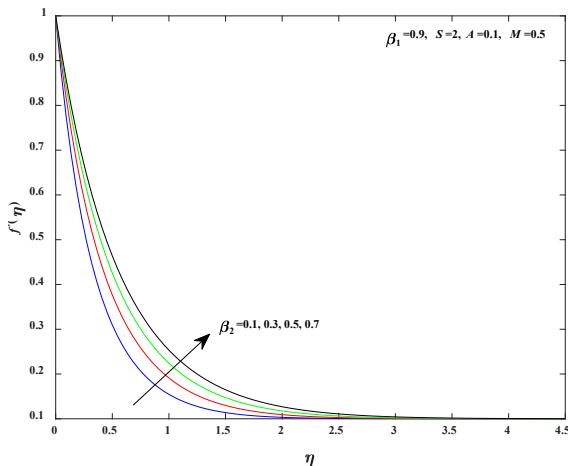
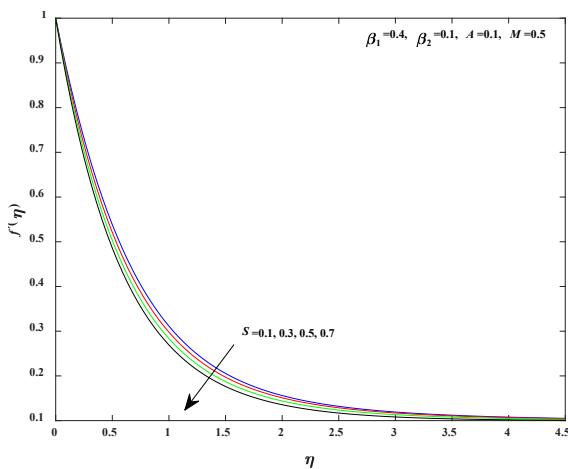
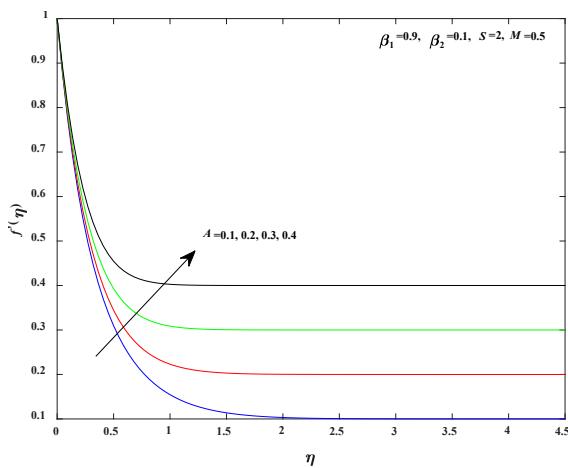
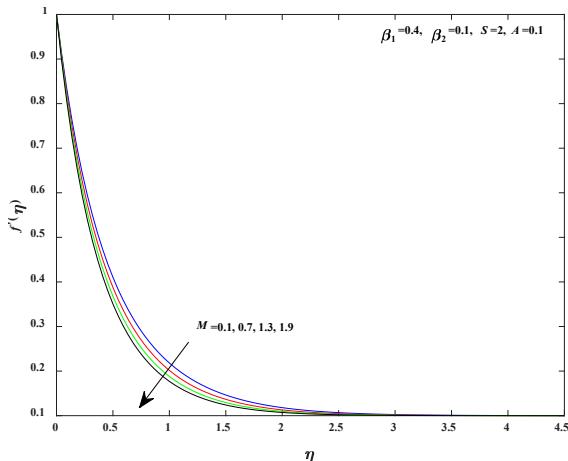


Figure 2. Influence of β_1 on the velocity field $f'(\eta)$

图 2. β_1 对速度的影响

**Figure 3.** Influence of β_2 on the velocity field $f'(\eta)$ **图 3.** β_2 对速度的影响**Figure 4.** Influence of S on the velocity field $f'(\eta)$ **图 4.** S 对速度的影响**Figure 5.** Influence of A on the velocity field $f'(\eta)$ **图 5.** A 对速度的影响

**Figure 6.** Influence of M on the velocity field $f'(\eta)$ **图 6. M 对速度的影响**

5. 总结

本文研究了磁场对拉伸板上 Oldroyd-B 流体非稳态驻点流动的影响。通过合适的相似变换将控制方程转化为常微分方程，并且使用双参数变形展开方法与基函数方法结合的方法(DPTEM-BF)进行求解。采用图形详细分析了各物性参数对速度的影响，值得注意的是：随着驻点参数的增大，流体的动量损失减小，边界层内流体速度越快趋于周围速度，所以速度边界层厚度变薄；而磁场参数的增加会使得磁场作用下的洛伦兹力增大，这个力阻碍了流体流动，因此流体流动变缓。

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