

# Generalized $\frac{\varphi'}{\varphi}$ -Expansion Method and the Traveling Wave Solutions of the STO Equation

Yuanyuan Han, Desheng Li, Ting Huang

School of Mathematics and System Science, Shenyang Normal University, Shenyang  
Email: [yuanyuan010529@163.com](mailto:yuanyuan010529@163.com)

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**Abstract:** This paper is about to discuss the method which is based on the generalized  $\frac{\varphi'}{\varphi}$ -expansion method, and explain how to determine  $\varphi$  by the equation itself without considering the auxiliary equation to make sure the solutions of the equation. This paper discusses the Sharma-Tasso-Olver (STO) equation and obtains the traveling wave solutions and trigonometric function solutions of the STO equation.

**Keywords:** Generalized  $\frac{\varphi'}{\varphi}$ -Expansion Method; Traveling Wave Solution; Trigonometric Functions Solution

## 扩展的 $\frac{\varphi'}{\varphi}$ -展开法及 Sharma-Tasso-Olver 方程的行波解

韩园媛, 李德生, 黄 婷

沈阳师范大学数学与系统科学学院, 沈阳  
Email: [yuanyuan010529@163.com](mailto:yuanyuan010529@163.com)

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**摘 要:** 本文在原有扩展的  $\frac{\varphi'}{\varphi}$ -展开法基础之上, 不考虑  $\varphi$  满足的辅助方程, 只利用方程本身来确定  $\varphi$ , 进而确定方程的解。文章讨论求解了 Sharma-Tasso-Olver(STO)方程, 获得了 STO 方程的行波解和三角函数解。

**关键词:** 扩展的  $\frac{\varphi'}{\varphi}$ -展开法; 行波解; 三角函数解

### 1. 引言

由于非线性偏微分方程的解对于解释物理学、生物学等领域出现的各种现象具有积极的意义, 因此非线性偏微分方程的求解受到了大家的广泛关注, 求解方法也层出不穷, 如 Darboux 变换<sup>[1]</sup>、Bäcklund 变换法<sup>[2-4]</sup>、辅助方程法<sup>[5]</sup>、Hirota 双线性方法<sup>[6]</sup>、对称法<sup>[7]</sup>等。

王明亮等人于 2008 年在文献[8]中提出了  $\frac{G'}{G}$  (即本文所说的  $\frac{\varphi'}{\varphi}$ )-展开法, 大量的文章利用这一方法求解非线性偏微分方程, 并对其进行了推广<sup>[9-15]</sup>。最近, 文献[16]的作者对这一方法进行了另一种推广, 但是仍然借助了辅助方程, 本文在此基础上, 去掉辅助方程这一条件, 只利用方程本身进行求解, 并利用该方法讨论研究了 STO 方程, 获得了该方程的精确解。

本文分为四个部分, 第一部分为引言; 第二部分详细描述了扩展的  $\frac{\varphi'}{\varphi}$ -展开法的步骤; 第三部利用这一方法讨论研究了 STO 方程, 获得了该方程的行波解和三角函数解; 第四部分阐述文章的结论。

## 2. 扩展的 $\frac{\varphi'}{\varphi}$ -展开法

利用扩展的  $\frac{\varphi'}{\varphi}$ -展开法求解非线性偏微分方程分为以下几个步骤:

- 1) 通过平衡方程的最高阶非线性项和最高阶导数项, 确定正整数  $n$ ;
- 2) 假设方程具有如下形式的解

$$u = A_0 + \sum_{k=1}^n \sum_{i+j=k} A_{ij} \left( \frac{\varphi_1'(\xi_1)}{\varphi_1(\xi_1)} \right)^i \left( \frac{\varphi_2'(\xi_2)}{\varphi_2(\xi_2)} \right)^j, \quad \sum_{i+j=k} A_{ij} \left( \frac{\varphi_1'(\xi_1)}{\varphi_1(\xi_1)} \right)^i \left( \frac{\varphi_2'(\xi_2)}{\varphi_2(\xi_2)} \right)^j \neq 0, \quad (1)$$

其中

$$\xi_1 = k_1 x + c_1 t + l_1, \quad \xi_2 = k_2 x + c_2 t + l_2, \quad \frac{\varphi_1'(\xi_1)}{\varphi_1(\xi_1)} = \frac{d\varphi_1(\xi_1)}{d\xi_1}, \quad \frac{\varphi_2'(\xi_2)}{\varphi_2(\xi_2)} = \frac{d\varphi_2(\xi_2)}{d\xi_2},$$

$A_0, A_{ij}, k_1, k_2, c_1, c_2, l_1, l_2$  为任意常数。将(1)式代入方程中, 由于  $\varphi_1^{-i} \varphi_2^{-j}$  ( $i = 0, 1, 2, \dots; j = 0, 1, 2, \dots$ ) 是线性无关的, 因此令  $\varphi_1^{-i} \varphi_2^{-j}$  的系数等于零, 得到关于  $A_0, A_{ij}, \varphi_1(\xi_1), \varphi_2(\xi_2)$  的常微分方程组;

- 3) 求解前一步获得的常微分方程组, 确定  $A_0, A_{ij}, \varphi_1(\xi_1), \varphi_2(\xi_2)$ ;
- 4) 利用前一步的结果以及(1)式, 即可得到方程的精确解。

## 3. 扩展的 $\frac{\varphi'}{\varphi}$ -展开法对 STO 方程的应用

考虑 Sharma-Tasso-Olver(STO)方程

$$u_t + 3\alpha u^2 u_x + \frac{3}{2} \alpha (u^2)_{xx} + \alpha u_{xxx} = 0, \quad (2)$$

其中  $\alpha$  是任意实常数。

平衡方程(2)的最高阶非线性项与最高阶导数项可得  $n = 1$ , 因此设方程(2)具有如下形式的解

$$u = A_0 + A_1 \frac{\varphi_1'(\xi_1)}{\varphi_1(\xi_1)} + A_2 \frac{\varphi_2'(\xi_2)}{\varphi_2(\xi_2)}, \quad (3)$$

其中  $\xi_1 = k_1 x + c_1 t + l_1, \xi_2 = k_2 x + c_2 t + l_2, A_0, A_1, A_2, k_1, k_2, c_1, c_2, l_1, l_2$  为任意常数, 将(3)式代入方程(2)中, 令  $\varphi_1^{-i} \varphi_2^{-j}$  的系数等于零, 可以得到关于  $A_0, A_1, A_2$  以及  $\varphi_1(\xi_1), \varphi_2(\xi_2)$  的常微分方程组, 求解一部分方程并整理其余的方程可得到下面的结果和方程

$$\begin{aligned} A_0 &= a, \quad A_1 = k_1, \quad A_2 = k_2, \\ c_1 \varphi_1' + 3\alpha a^2 k_1 \varphi_1' + 3\alpha a k_1^2 \varphi_1'' + \alpha k_1^3 \varphi_1''' &= 0, \end{aligned} \quad (4)$$

$$c_2\varphi_2' + 3\alpha a^2 k_2\varphi_2' + 3\alpha a k_2^2\varphi_2'' + \alpha k_2^3\varphi_2''' = 0, \quad (5)$$

$$k_1\varphi_1''\varphi_2' + 2a\varphi_1'\varphi_2' + k_2\varphi_1'\varphi_2'' = 0, \quad (6)$$

其中  $a$  为任意常数。

方程(4)的特征方程为

$$\lambda \left[ (c_1 + 3\alpha a^2 k_1) + 3\alpha a k_1^2 \lambda + \alpha k_1^3 \lambda^2 \right] = 0,$$

记方程

$$(c_1 + 3\alpha a^2 k_1) + 3\alpha a k_1^2 \lambda + \alpha k_1^3 \lambda^2 = 0$$

的判别式和根分别为

$$\Delta_1 = -3\alpha^2 a^2 k_1^4 - 4\alpha c_1 k_1^3, \quad \lambda_{11} = \frac{-3\alpha a k_1^2 + \sqrt{\Delta_1}}{2\alpha k_1^3}, \quad \lambda_{12} = \frac{-3\alpha a k_1^2 - \sqrt{\Delta_1}}{2\alpha k_1^3},$$

同理，对于方程(5)的特征方程中的判别式和特征根记为

$$\Delta_2 = -3\alpha^2 a^2 k_2^4 - 4\alpha c_2 k_2^3, \quad \lambda_{21} = \frac{-3\alpha a k_2^2 + \sqrt{\Delta_2}}{2\alpha k_2^3}, \quad \lambda_{22} = \frac{-3\alpha a k_2^2 - \sqrt{\Delta_2}}{2\alpha k_2^3}.$$

$\Delta_1, \Delta_2$  的符号决定了方程(4)和方程(5)的解的形式，进而决定了方程(2)的解的形式，按解的形式不同，可以分为以下三种情况：

$$\text{情况一} \quad \begin{cases} \Delta_1 > 0, \\ \Delta_2 > 0, \end{cases}$$

1) 显然

$$\varphi_1(\xi_1) = M e^{\lambda_{11}\xi_1} + b_1(t), \quad (7)$$

$$\varphi_2(\xi_2) = N e^{\lambda_{21}\xi_2} + b_2(t), \quad (8)$$

分别为方程(4)和方程(5)的解，其中  $M, N$  是任意常数， $b_1(t), b_2(t)$  是  $t$  的任意函数。将其代入方程(6)

$$2a + k_1\lambda_{11} + k_2\lambda_{21} = 0, \quad (9)$$

事实上，(9)式为  $k_1, k_2, c_1, c_2, a, \alpha$  满足的方程，为了书写方便，我们用(9)式表示。将(7)、(8)式代入(3)式，并利用(9)式，可得方程(2)的行波解

$$u_{11} = \frac{k_1\lambda_{11}}{2} \tanh\left(\frac{\lambda_{11}}{2}\xi_1 + \xi_{10}\right) + \frac{k_2\lambda_{21}}{2} \tanh\left(\frac{\lambda_{21}}{2}\xi_2 + \xi_{20}\right), \quad (10)$$

其中

$$\xi_{10} = \frac{1}{2} \ln \frac{M}{b_1(t)}, \quad \xi_{20} = \frac{1}{2} \ln \frac{N}{b_2(t)}, \quad \xi_1 = k_1 x + c_1 t + l_1, \quad \xi_2 = k_2 x + c_2 t + l_2, \quad k_1, k_2, c_1, c_2, a, \alpha$$

满足(9)式。

2) 显然

$$\varphi_1(\xi_1) = M e^{\lambda_{11}\xi_1} + b_1(t), \quad (11)$$

$$\varphi_2(\xi_2) = N e^{\lambda_{22}\xi_2} + b_2(t), \quad (12)$$

分别为方程(4)和方程(5)的解, 由(11)、(12)式可以得到方程(2)的行波解为

$$u_{12} = \frac{k_1 \lambda_{11}}{2} \tanh\left(\frac{\lambda_{11}}{2} \xi_1 + \xi_{10}\right) + \frac{k_2 \lambda_{22}}{2} \tanh\left(\frac{\lambda_{22}}{2} \xi_2 + \xi_{20}\right), \quad (13)$$

其中

$$\xi_{10} = \frac{1}{2} \ln \frac{M}{b_1(t)}, \quad \xi_{20} = \frac{1}{2} \ln \frac{N}{b_2(t)}, \quad \xi_1 = k_1 x + c_1 t + l_1, \quad \xi_2 = k_2 x + c_2 t + l_2, \quad k_1, k_2, c_1, c_2, a, \alpha$$

满足如下方程

$$2a + k_1 \lambda_{11} + k_2 \lambda_{22} = 0 \quad (14)$$

3) 显然

$$\varphi_1(\xi_1) = M e^{\lambda_{12} \xi_1} + b_1(t), \quad (15)$$

$$\varphi_2(\xi_2) = N e^{\lambda_{22} \xi_2} + b_2(t), \quad (16)$$

分别为方程(4)和方程(5)的解, 由(15)、(16)式可以得到方程(2)的行波解为

$$u_{13} = \frac{k_1 \lambda_{12}}{2} \tanh\left(\frac{\lambda_{12}}{2} \xi_1 + \xi_{10}\right) + \frac{k_2 \lambda_{22}}{2} \tanh\left(\frac{\lambda_{22}}{2} \xi_2 + \xi_{20}\right) \quad (17)$$

其中

$$\xi_{10} = \frac{1}{2} \ln \frac{M}{b_1(t)}, \quad \xi_{20} = \frac{1}{2} \ln \frac{N}{b_2(t)}, \quad \xi_1 = k_1 x + c_1 t + l_1, \quad \xi_2 = k_2 x + c_2 t + l_2, \quad k_1, k_2, c_1, c_2, a, \alpha$$

满足如下方程

$$2a + k_1 \lambda_{12} + k_2 \lambda_{22} = 0. \quad (18)$$

$$\text{情况二} \quad \begin{cases} \Delta_1 > 0, \\ \Delta_2 < 0, \end{cases}$$

1) 显然

$$\varphi_1(\xi_1) = M e^{\lambda_{11} \xi_1} + b_1(t), \quad (19)$$

$$\varphi_2(\xi_2) = N_1 e^{\frac{3a}{2k_2} \cos\left(\frac{\sqrt{-\Delta_2}}{2\alpha k_2^3} \xi_2\right)} + N_2 e^{\frac{3a}{2k_2} \sin\left(\frac{\sqrt{-\Delta_2}}{2\alpha k_2^3} \xi_2\right)} + b_2(t), \quad (20)$$

分别为方程(4)和方程(5)的解, 其中  $M, N_1, N_2$  为任意常数,  $b_1(t), b_2(t)$  是  $t$  的任意函数, 将其代入(6)式, 并考虑  $\Delta_2 = -3\alpha^2 a^2 k_2^4 - 4\alpha c_2 k_2^3$  可得

$$c_2 = -3\alpha a^2 k_2, \quad (21)$$

将(19)~(21)式代入(3)式可得方程(2)的精确解

$$u_{21} = a + \frac{k_1 \lambda_{11}}{2} + \frac{k_1 \lambda_{11}}{2} \tanh\left(\frac{\lambda_{11}}{2} \xi_1 + \xi_{10}\right) + k_2 \cdot \frac{\varphi_2'}{\varphi_2}, \quad (22)$$

其中

$$\xi_{10} = \frac{1}{2} \ln \frac{M}{b_1(t)}, \quad \xi_1 = k_1 x + c_1 t + l_1, \quad \xi_2 = k_2 x - 3\alpha a^2 k_2 t + l_2.$$

2) 显然

$$\varphi_1(\xi_1) = M e^{\lambda_{12}\xi_1} + b_1(t), \quad (23)$$

$$\varphi_2(\xi_2) = N_1 e^{-\frac{3a}{2k_2}\xi_2} \cos\left(\frac{\sqrt{-\Delta_2}}{2\alpha k_2^3}\xi_2\right) + N_2 e^{-\frac{3a}{2k_2}\xi_2} \sin\left(\frac{\sqrt{-\Delta_2}}{2\alpha k_2^3}\xi_2\right) + b_2(t), \quad (24)$$

分别为方程(4)和方程(5)的解, 由(23)、(24)式可以得到方程(2)的解为

$$u_{22} = a + \frac{k_1\lambda_{12}}{2} + \frac{k_1\lambda_{12}}{2} \tanh\left(\frac{\lambda_{12}}{2}\xi_1 + \xi_{10}\right) + k_2 \cdot \frac{\varphi_2'}{\varphi_2}, \quad (25)$$

其中

$$\xi_{10} = \frac{1}{2} \ln \frac{M}{b_1(t)}, \quad \xi_1 = k_1 x + c_1 t + l_1, \quad \xi_2 = k_2 x - 3\alpha a^2 k_2 t + l_2.$$

情况三  $\begin{cases} \Delta_1 < 0, \\ \Delta_2 < 0, \end{cases}$

显然

$$\varphi_1 = M_1 e^{-\frac{3a}{2k_1}\xi_1} \cos\left(\frac{\sqrt{-\Delta_1}}{2\alpha k_1^3}\xi_1\right) + M_2 e^{-\frac{3a}{2k_1}\xi_1} \sin\left(\frac{\sqrt{-\Delta_1}}{2\alpha k_1^3}\xi_1\right) + b_1(t), \quad (26)$$

$$\varphi_2 = N_1 e^{-\frac{3a}{2k_2}\xi_2} \cos\left(\frac{\sqrt{-\Delta_2}}{2\alpha k_2^3}\xi_2\right) + N_2 e^{-\frac{3a}{2k_2}\xi_2} \sin\left(\frac{\sqrt{-\Delta_2}}{2\alpha k_2^3}\xi_2\right) + b_2(t), \quad (27)$$

分别为方程(4)和方程(5)的解, 其中  $M_1, M_2, N_1, N_2$  为任意常数,  $b_1(t), b_2(t)$  为  $t$  的任意函数, 将其代入(6)式有

$$-AM_1 N_1 + BM_1 N_2 + CM_2 N_1 - DM_2 N_2 = 0, \quad (28)$$

$$-BM_1 N_1 - AM_1 N_2 + DM_2 N_1 + CM_2 N_2 = 0, \quad (29)$$

$$-CM_1 N_1 + DM_1 N_2 - AM_2 N_1 + BM_2 N_2 = 0, \quad (30)$$

$$DM_1 N_1 + CM_1 N_2 + BM_2 N_1 + AM_2 N_2 = 0, \quad (31)$$

其中

$$A = \frac{3a}{8k_1 k_2} \left( 6a^2 + \frac{\Delta_1}{\alpha^2 k_1^4} + \frac{\Delta_2}{\alpha^2 k_2^4} \right), \quad B = \frac{\sqrt{-\Delta_2}}{8\alpha k_1 k_2^3} \left( 15a^2 + \frac{\Delta_1}{\alpha^2 k_1^4} \right),$$

$$C = \frac{\sqrt{-\Delta_1}}{8\alpha k_1^3 k_2} \left( 15a^2 + \frac{\Delta_2}{\alpha^2 k_2^4} \right), \quad D = \frac{a\sqrt{\Delta_1 \Delta_2}}{\alpha^2 k_1^3 k_2^3}.$$

因此方程(2)的精确解为

$$u_3 = a + k_1 \cdot \frac{\varphi_1'}{\varphi_1} + k_2 \cdot \frac{\varphi_2'}{\varphi_2}, \quad (32)$$

其中  $\varphi_1, \varphi_2$  分别为(26)、(27)式,  $k_1, k_2, c_1, c_2, a, \alpha, M_1, M_2, N_1, N_2$  满足(28)~(31)式。

事实上, 根据  $\Delta_1, \Delta_2$  的符号不同进行组合, 一共有 9 种情况, 但是  $\begin{cases} \Delta_1 > 0, \\ \Delta_2 = 0, \end{cases} \begin{cases} \Delta_1 > 0, \\ \Delta_2 > 0, \end{cases} \begin{cases} \Delta_1 = 0, \\ \Delta_2 > 0, \end{cases}$  以及  $\begin{cases} \Delta_1 = 0, \\ \Delta_2 = 0, \end{cases}$  获

得的解的形式与情况一中解的形式相同,  $\begin{cases} \Delta_1 < 0, \\ \Delta_2 > 0, \end{cases} \begin{cases} \Delta_1 = 0, \\ \Delta_2 < 0, \end{cases}$  以及  $\begin{cases} \Delta_1 < 0, \\ \Delta_2 = 0, \end{cases}$  获得的解的形式与情况二中解的形式相

同, 这里不再赘述。

#### 4. 结论

实际上, 解  $u_{11}, u_{12}, u_{13}$  是方程(2)的双孤子解, 由于  $\lambda_{11}, \lambda_{12}$  异号,  $\lambda_{21}, \lambda_{22}$  异号, 使得表达式中  $t$  的系数不同, 即波速不同, 所以  $u_{11}, u_{12}, u_{13}$  是三个不同的双孤子解; 同理,  $u_{21}, u_{22}$  是两个不同的复合形式解, 这种复合形式的解由双曲正切函数和三角函数组成;  $u_3$  是包含多个任意参数的三角函数解。本文获得的解中都包含多个任意参数, 若将参数选为某特殊值, 将会得到文献[16]中的部分解。

在解  $u_{11}$  中, 选取

$$\begin{aligned} a &= 0, \quad k_1 = F_1, \quad k_2 = F_2, \quad l_1 = F_3, \quad l_2 = F_4, \\ c_1 &= -\alpha F_1^3 (\lambda_1^2 - 4\mu_1), \quad c_2 = -\alpha F_2^3 (\lambda_2^2 - 4\mu_2), \\ \xi_{10} &= \frac{1}{2} \ln \frac{M}{b_1(t)} = \arctan \frac{M_1}{M_2}, \quad \xi_{20} = \frac{1}{2} \ln \frac{N}{b_2(t)} = \arctan \frac{N_1}{N_2}, \end{aligned} \quad (33)$$

并利用(9)式, 则(10)式变为

$$\begin{aligned} u_{11} &= \frac{F_1 \sqrt{\lambda_1^2 - 4\mu_1}}{2} \tanh \left( \frac{\sqrt{\lambda_1^2 - 4\mu_1}}{2} \xi_1 + \xi_{10} \right) + \frac{F_2 \sqrt{\lambda_2^2 - 4\mu_2}}{2} \tanh \left( \frac{\sqrt{\lambda_2^2 - 4\mu_2}}{2} \xi_2 + \xi_{20} \right), \\ \xi_1 &= F_1 x - \frac{\alpha F_1 [F_1^2 (\lambda_1^2 - 4\mu_1) + 3F_2^2 (\lambda_2^2 - 4\mu_2)] t}{4} + F_3, \quad \xi_{10} = \arctan \frac{M_1}{M_2}, \\ \xi_2 &= F_2 x - \frac{\alpha F_2 [F_2^2 (\lambda_2^2 - 4\mu_2) + 3F_1^2 (\lambda_1^2 - 4\mu_1)] t}{4} + F_4, \quad \xi_{20} = \arctan \frac{N_1}{N_2}, \end{aligned} \quad (34)$$

(34)式即为文献[16]中的(13)式。

在解  $u_3$  中, 选取

$$\begin{aligned} a &= 0, \quad k_1 = F_1, \quad k_2 = F_2, \quad l_1 = F_3, \quad l_2 = F_4, \\ c_1 &= -\frac{\alpha F_1^3 (\lambda_1^2 - 4\mu_1)}{4}, \quad c_2 = -\frac{\alpha F_2^3 (\lambda_2^2 - 4\mu_2)}{4}, \end{aligned}$$

交换  $M_1, M_2$  的位置,  $N_1, N_2$  的位置, 并利用(28)~(31)式, 则(32)式变为

$$\begin{aligned} u_3 &= \frac{F_1 \sqrt{-(\lambda_1^2 - 4\mu_1)}}{2} \cdot \frac{M_1 \cos \left( \frac{\sqrt{-(\lambda_1^2 - 4\mu_1)}}{2} \xi_1 \right) - M_2 \sin \left( \frac{\sqrt{-(\lambda_1^2 - 4\mu_1)}}{2} \xi_1 \right)}{M_1 \sin \left( \frac{\sqrt{-(\lambda_1^2 - 4\mu_1)}}{2} \xi_1 \right) + M_2 \cos \left( \frac{\sqrt{-(\lambda_1^2 - 4\mu_1)}}{2} \xi_1 \right)}, \\ &+ \frac{F_2 \sqrt{-(\lambda_2^2 - 4\mu_2)}}{2} \cdot \frac{N_1 \cos \left( \frac{\sqrt{-(\lambda_2^2 - 4\mu_2)}}{2} \xi_2 \right) - N_2 \sin \left( \frac{\sqrt{-(\lambda_2^2 - 4\mu_2)}}{2} \xi_2 \right)}{N_1 \sin \left( \frac{\sqrt{-(\lambda_2^2 - 4\mu_2)}}{2} \xi_2 \right) + N_2 \cos \left( \frac{\sqrt{-(\lambda_2^2 - 4\mu_2)}}{2} \xi_2 \right)}, \end{aligned}$$

$$\begin{aligned}\xi_1 &= F_1 x - \frac{\alpha F_1 \left[ F_1^2 (\lambda_1^2 - 4\mu_1) + 3F_2^2 (\lambda_2^2 - 4\mu_2) \right] t}{16} + F_3, \\ \xi_2 &= F_2 x - \frac{\alpha F_2 \left[ F_2^2 (\lambda_2^2 - 4\mu_2) + 3F_1^2 (\lambda_1^2 - 4\mu_1) \right] t}{16} + F_4,\end{aligned}\tag{35}$$

(35)式即为文献[16]中的(14)式。

利用本文的方法可以获得很多非线性偏微分方程的双孤子解、三角函数解，且都是行波解，能否利用该方法得到非行波解还是一个有待解决的问题，有兴趣的读者可以做进一步的研究。

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